

to  $\Delta S^0$  of the  $\text{CN}^-$ . Thus, the observed  $\Delta S^0$  is essentially that resulting from the contribution of the proton. The molecules previously studied by

Canady, *et al.*, and by us had low polarity and, thus, did not show this effect to an appreciable extent.

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## The Electronic Structures and Magnetic Properties of the Chromyl and Molybdenyl Ions

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The electron spin resonance spectra and the magnetic susceptibilities of  $(\text{NH}_4)_2[\text{MoOCl}_5]$  and  $\text{Rb}_2[\text{CrOCl}_5]$  are reported. The results are consistent with an electronic structural description which includes considerable oxygen-to-metal  $\pi$ -bonding. Complete consistency between the e.s.r. and the susceptibility data is found for  $(\text{NH}_4)_2[\text{MoOCl}_5]$ . The lack of consistency of the data for  $\text{Rb}_2[\text{CrOCl}_5]$  is attributed to the presence of a small amount of  $\text{Cr}^{3+}$  impurity. Comparison of e.s.r. data with reported susceptibility data for  $\text{Mo}^{6+}$  solutions of varying HCl concentration is interpreted in terms of the nature of the species present.

### Introduction

In a previous report<sup>1</sup> the electronic structures and spectra of the chromyl ( $\text{CrO}^{3+}$ ) and molybdenyl ( $\text{MoO}^{3+}$ ) ions were discussed. Both the charge transfer and the "crystal field" d-d spectra could be accounted for by the molecular orbital energy level scheme derived for the electronically equivalent ( $d^1$ ) vanadyl ( $\text{VO}^{2+}$ ) ion.<sup>2</sup> The most striking feature of the electronic structures of these oxycations is the existence of considerable oxygen-to-metal  $\pi$ -bonding.

In the present paper the electron spin resonance spectra and the temperature dependencies of the magnetic susceptibilities of  $(\text{NH}_4)_2[\text{MoOCl}_5]$  and  $\text{Rb}_2[\text{CrOCl}_5]$  are investigated and the results are interpreted in terms of the previously proposed electronic structural description. The magnetic susceptibilities of solid samples of  $(\text{NH}_4)_2[\text{MoOCl}_5]$ <sup>3</sup> and  $\text{Rb}_2[\text{CrOCl}_5]$ <sup>4</sup> and solutions of  $\text{Mo}^{5+}$ <sup>5</sup> have been reported in the literature, but for only room temperature.

A comparison of the results of the present e.s.r. measurements on  $\text{Mo}^{5+}$  solutions with the susceptibility results of Sacconi and Cini<sup>5</sup> allows an

interpretation of the magnetic nature of the species present in these solutions.

### Results

**Chromyl Ion.**—The electron spin resonance spectrum of a solid sample of  $\text{Rb}_2[\text{CrOCl}_5]$  shows a single broad resonance line centered at  $g = 1.986$ . The  $g$  value indicates an orbitally non-degenerate ground state, as expected for a compound containing  $\text{CrO}^{3+}$ . The  $g$  factor for a paramagnetic molecule with an orbitally non-degenerate ground state differs from 2.0023 due to the mixing of ground state and excited states *via* spin-orbit coupling. The expressions for the  $g$  values for  $d^1$  transition metal oxycations in  $C_{4v}$  ligand fields are<sup>2</sup>

$$g_{\perp} = 2.0023 \left( 1 - \frac{\xi}{\Delta E(^2B_2 \rightarrow ^2E(I))} \right) \quad (1)$$

$$g_{\parallel} = 2.0023 \left( 1 - \frac{4\xi}{\Delta E(^2B_2 \rightarrow ^2B_1)} \right) \quad (2)$$

with

$$\langle g \rangle^2 = 1/3 (2g_{\perp}^2 + g_{\parallel}^2) \quad (3)$$

In these equations  $\xi$  is the spin-orbit coupling constant for  $\text{Cr}^{5+}$  in  $\text{CrO}^{3+}$  and the  $\Delta E$ 's are the energies of known crystal field transitions in  $\text{Rb}_2[\text{CrOCl}_5]$ .<sup>1</sup> The  $\xi$  value for  $\text{CrO}^{3+}$  is calculated to be  $65 \text{ cm.}^{-1}$ , using the observed  $\langle g \rangle$  value and the known  $\Delta E$  values ( $\Delta E(^2B_2 \rightarrow ^2E(I)) = 12,900 \text{ cm.}^{-1}$ ,  $\Delta E(^2B_2 \rightarrow ^2B_1) = 23,500 \text{ cm.}^{-1}$ ).

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The predicted values of  $g_{\perp}$  and  $g_{\parallel}$  are, respectively, 1.990 and 1.978.

The magnetic susceptibility of  $\text{Rb}_2[\text{CrOCl}_5]$  was measured over the range 77–318°K. These diamagnetic corrected data are given in Table I. The theoretical expression for magnetic susceptibility given by Van Vleck<sup>6</sup> is

$$\chi = (C/T) + \chi_{\text{H.F.}} \quad (4)$$

where  $C$  is the Curie constant and  $\chi_{\text{H.F.}}$  is the temperature-independent contribution to the susceptibility. In this case

$$C = N\beta^2\langle g \rangle^2/4k \quad (5)$$

where  $N$  is Avogadro's Number,  $\beta$  is the Bohr magneton,  $k$  is Boltzmann's constant, and

$$\chi_{\text{H.F.}} = \frac{2N\beta^2}{g_m} \left[ \frac{2}{3} \sum_i \frac{|\langle \psi_0 | L_x + 2S_x | \psi_i \rangle|^2}{\Delta E_{0,i}} + \frac{1}{3} \sum_i \frac{|\psi_0 L_x + 2S_x \psi_i|^2}{\Delta E_{0,i}} \right] \quad (6)$$

where  $\Delta E_{0,i}$  represents the energy differences from the ground state  $\psi_0$  to excited states  $\psi_i$ .

For  $\text{Rb}_2[\text{CrOCl}_5]$ , with  $\langle g \rangle = 1.986$  and  $\Delta E'$  from the spectral data, eq. 4 gives

$$\chi = (0.370/T) + 55 \times 10^{-6} \quad (7)$$

A plot of  $\chi$  vs.  $(1/T)$  for  $\text{Rb}_2[\text{CrOCl}_5]$  is shown in Fig. 1. This plot gives

$$\chi = 0.536/T \quad (8)$$

in very poor agreement with the calculated equation. The compound is unstable, and even at room temperature decomposition to give  $\text{Cr}^{3+}$  and  $\text{Cr}^{6+}$  takes place; total decomposition occurs above 45°. While care was taken to reduce handling times and to protect the samples from oxidative conditions and moisture, it is probable that some decomposition did take place. A rough correction was applied to the data to account for the susceptibility of  $\text{Cr}^{3+}$  impurity. It was found that a 7% impurity of  $\text{Cr}^{3+}$  would account for the deviations from the expected behavior. This correction is only possible because the effective magnetic moment of the impurity does not vary with temperature.<sup>7</sup> Further refinement of the data is not justified.

**Molybdenyl Ion.**—Detailed spectral studies<sup>1,8,9</sup> have shown that the ion  $\text{MoOCl}_5^{2-}$  is present as a monomer in 10–12 *M* HCl solutions of  $(\text{NH}_4)_2$ -

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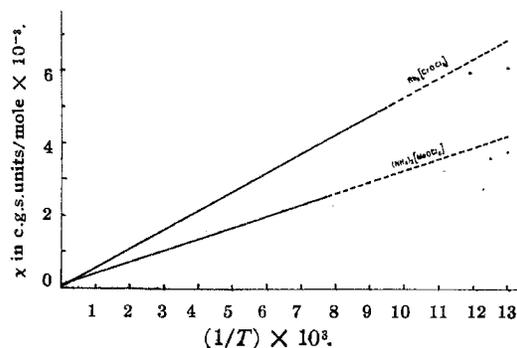


Fig. 1.—Magnetic susceptibilities of  $\text{Rb}_2[\text{CrOCl}_5]$  and  $(\text{NH}_4)_2[\text{MoOCl}_5]$  vs.  $1/T$ . All experimental points fall on the lines except as indicated.

TABLE I  
MAGNETIC SUSCEPTIBILITIES OF  $(\text{NH}_4)_2[\text{MoOCl}_5]$  AND  $\text{Rb}_2[\text{CrOCl}_5]$

$(\text{NH}_4)_2[\text{MoOCl}_5]$ diamagnetic correction, $153.9 \times 10^{-6}$		$\text{Rb}_2[\text{CrOCl}_5]$ diamagnetic correction, $165.3 \times 10^{-6}$	
°K.	$\chi'_M^a$	°K.	$\chi'_M^a$
77	3742	77	6065
80	3619	84	5930
129	2519	145	3741
189	1777	152	3524
195	1754	159	3420
201	1706	164	3258
206	1654	170	3141
212	1605	177	3046
217	1556	183	2958
223	1530	189	2792
228	1514	195	2738
233	1484	201	2650
238	1455	206	2571
243	1416	212	2517
248	1374	217	2442
255	1348	223	2375
261	1325	228	2309
267	1299	233	2246
273	1273	238	2209
295	1195	243	2192
306	1169	248	2142
333	1075	255	2113
375	958	261	2030
424	853	267	1984
471	766	273	1934
		294	1797
		301	1797
		308	1772
		318	1726

<sup>a</sup>  $\chi'_M$  is the molar susceptibility in units  $\times 10^{-6}$  c.g.s./mole corrected for diamagnetism.

$[\text{MoOCl}_5]$ . The electron spin resonance spectrum of  $\text{MoO}^{3+}$  in 10–12 *M* HCl shows a strong central line centered at  $g = 1.947$ . The central line is flanked by three weak satellites on each side. The strong central line is due to the  $^{96}\text{Mo}$  ( $I = 0$ ; abundance, 74.82%) isotope; the weaker

satellites are due to the  $^{96}\text{Mo}$  ( $I = 5/2$ ; abundance, 15.78%) and  $^{97}\text{Mo}$  ( $I = 5/2$ ; abundance, 9.6%) isotopes. The latter two isotopes have approximately the same nuclear moment and thus the six weak satellites were not resolved into hyperfine contributions due to the individual isotopes. The average hyperfine splitting due to the  $^{96}\text{Mo}$  and  $^{97}\text{Mo}$  isotopes is about 50 gauss.

The intensity of the e.s.r. spectrum of a  $\text{MoO}^{3+}$  solution decreases drastically going from 10 to 4  $M$  HCl; at 4  $M$  HCl the e.s.r. signal is barely detectable. It is also important to note that the width of the central line stays constant at about 8 gauss, independent of the HCl concentration. The percentage of  $\text{MoO}^{3+}$  monomer present in solution at the various HCl concentrations is directly proportional to the intensity of the e.s.r. spectrum. The analysis of these e.s.r. intensities in terms of the percentage of  $\text{MoO}^{3+}$  monomer present is given in Table II.

TABLE II

Concentration of HCl, $M$	% total susceptibility <sup>a</sup>	% $\text{MoO}^{3+}$ monomer from e.s.r. results	Difference
12	100	100	0
10	100	100	0
8.4	99	85	14
7.1	95	74	21
6	75	46	29
5.2	41	16	25
4	9	about 1	8
2	0	0	0

<sup>a</sup> From the data of Sacconi and Cini (see ref. 5).

For the  $\text{MoOCl}_5^{2-}$  monomer, with  $\langle g \rangle = 1.947$ ,  $\Delta E(^2B_2 \rightarrow ^2E(I)) = 13,800 \text{ cm.}^{-1}$ , and  $\Delta E(^2B_2 \rightarrow ^2B_1) = 23,000 \text{ cm.}^{-1}$ , the spin-orbit constant,  $g_{\parallel}$ , and  $g_{\perp}$  can be obtained from eq. 1, 2, and 3. This calculation gives  $g_{\perp} = 1.965$ ,  $g_{\parallel} = 1.915$ , and  $\xi = 240 \text{ cm.}^{-1}$ .

The magnetic susceptibility of a solid sample of  $(\text{NH}_4)_2[\text{MoOCl}_5]$  was measured between 77 and 470°K. These results are given in Table I. For  $(\text{NH}_4)_2[\text{MoOCl}_5]$ , with  $\langle g \rangle = 1.947$  and the  $\Delta E$ 's from the spectral data,<sup>1</sup> eq. 4 becomes

$$\chi = (0.356/T) + 55 \times 10^{-6} \quad (9)$$

A plot of  $\chi$  vs.  $1/T$  is shown in Fig. 2. The linear portion of this curve gives  $C = 0.333$  and  $\chi_{\text{H.F.}} = 90 \times 10^{-6}$ , in satisfactory agreement with the predicted equation (9).

Equation 9 does not account for the "tailing off" of the  $\chi$  vs.  $1/T$  curve at low temperatures. This non-linearity may be accounted for by

including the Weiss constant in the temperature dependent term of eq. 9. This more detailed analysis of the experimental susceptibility results gives the expression

$$\chi = 0.35/(T + 15) + 75 \times 10^{-6} \quad (10)$$

which is in excellent agreement with theory. The average effective moment corrected for the temperature independent paramagnetism is found to be 1.67 B.M., which of course is expected, since  $\mu^2 = 3/4 \langle g \rangle^2$ . The small  $\theta$  value of  $+15^\circ\text{K.}$  may be due to a small degree of antiferromagnetism ( $J = 4^\circ\text{K.}$ ) or a lack of magnetic dilution.

### Discussion

The electronic structural similarities of  $\text{VO}^{2+}$ ,  $\text{CrO}^{3+}$ , and  $\text{MoO}^{3+}$ , concluded from detailed spectral studies, already have been elaborated.<sup>1</sup> The existence of considerable oxygen-to-metal  $\pi$ -bonding seems to be the principal feature of the electronic structure of these  $d^1$  oxycations.

The magnetic properties of  $(\text{NH}_4)_2[\text{MoOCl}_5]$  and  $\text{Rb}_2[\text{CrOCl}_5]$  are completely consistent with this electronic structural description. The  $\langle g \rangle$  values are very nearly 2, indicating that the orbital angular momentum is effectively quenched in the low symmetry ligand field. The free ion spin-orbit coupling constants for  $\text{Cr}^{5+}$  ( $380 \text{ cm.}^{-1}$ ) and  $\text{Mo}^{5+}$  ( $1030 \text{ cm.}^{-1}$ )<sup>10</sup> are drastically reduced in the oxycations; for  $\text{CrO}^{3+}$ ,  $\xi = 65 \text{ cm.}^{-1}$ ; for  $\text{MoO}^{3+}$ ,  $\xi = 240 \text{ cm.}^{-1}$ . This reduction in  $\xi$  also is observed for  $\text{VO}^{2+}$  and is accounted for by the oxygen-to-metal  $\pi$ -bonding.<sup>2</sup>

The structures of  $\text{Mo}^{5+}$  species in solutions of varying HCl concentration deserve special attention. It may be safely concluded from previous spectral studies and the present e.s.r. investigation that the monomer  $\text{MoO}^{3+}$  is the only species present in 10–12  $M$  HCl solutions. The situation is more complicated in solutions less than 10  $M$  in HCl. There are now three different types of information available on these solutions; (1) the magnetic susceptibility results reported by Sacconi and Cini,<sup>5</sup> (2) the electronic spectral results reported by Haight,<sup>9</sup> and (3) the present e.s.r. investigation. The bulk magnetic susceptibility measurements represent all the paramagnetic species in solution. The e.s.r. experiments are able to give just the amount of  $\text{MoO}^{3+}$  monomer present at any given HCl concentration. Table II affords a comparison of the magnetic



sample sealed in a 3-mm. tube under nitrogen. The solutions used for the measurement of the spectrum of  $(\text{NH}_4)_2[\text{MoOCl}_6]$  were 0.03 *M*. It was noted that the line width of this compound does not change on dilution to 0.003 *M*, therefore the more concentrated solution was chosen to obtain a stronger signal. The strength of a given line was measured by the area under the curve. All e.s.r. measurements were made at  $19 \pm 1^\circ$ . The magnetic susceptibility measurements of Sacconi and Cini<sup>6</sup> on  $\text{MoO}^{3+}$  solutions were made at  $20^\circ$ .

(3) **Magnetic Susceptibilities.**—The magnetic susceptibility measurements were carried out on a Gouy balance using an electromagnet with a maximum field strength of about 3000 gauss. The balance was fitted with a suitable

dewar flask to contain the cooling liquids. Temperatures between 77 and 273°K. were achieved by using precooled isopentane and allowing it to warm up slowly.

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## The Structure of the Diammoniate of Dimethylgallium Chloride, $(\text{CH}_3)_2\text{ClGa}\cdot 2\text{NH}_3$ , and of Ethylenediamine-Dimethylgallium Chloride

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Chemical evidence supports the structure  $[(\text{CH}_3)_2\text{Ga}(\text{NH}_3)_2]\text{Cl}$  for the known compound  $(\text{CH}_3)_2\text{GaCl}\cdot 2\text{NH}_3$ , reported by Kraus and Toonder and the structure  $[(\text{CH}_3)_2\text{Ga}(\text{H}_2\text{NCH}_2\text{CH}_2\text{NH}_2)]\text{Cl}$  for the new ethylenediamine analog. The latter compound is stable in water for limited periods of time.

The compound  $(\text{CH}_3)_2\text{GaCl}\cdot 2\text{NH}_3$  was reported by Kraus and Toonder<sup>1</sup> in 1933, but the structure was never established. An ammonium model, analogous to the early ammonia model used for the diammoniate of diborane, was suggested by Coates<sup>2</sup>:  $[\text{NH}_4^+][(\text{CH}_3)_2\text{GaNH}_2\text{Cl}^-]$ ; however, little real structural evidence was ever available. Following recognition of the formula  $[\text{H}_2\text{B}(\text{NH}_3)_2][\text{BH}_4^-]$  for the diammoniate of diborane,<sup>3</sup> another possible structure for the diammoniate of dimethylgallium chloride appeared probable:  $[(\text{CH}_3)_2\text{Ga}(\text{NH}_3)_2]\text{Cl}$ . Chemical evidence supports this latter model.

### Experimental

**General Procedures.**—Standard vacuum line techniques were used with moisture- and oxygen-sensitive materials.<sup>4</sup> Filtration of air-sensitive materials was per-

formed with the vacuum line filtration assembly which is described elsewhere.<sup>5</sup>

X-Ray powder patterns were obtained in 5.7 and 11.4 cm. diameter cameras with  $\text{CuK}\alpha$  radiation. Films were read on a North American Phillips film reader and intensities were estimated visually. Infrared spectra were obtained on a Perkin-Elmer Model 21 double beam recording spectrophotometer with NaCl optics; KCl disk and Nujol mull techniques were employed. Calibration of the spectrometer was checked with indene.<sup>6</sup>

**Materials.** 1.  $\text{Ga}(\text{CH}_3)_3$  was prepared by the action of 67.8 g. of  $\text{Hg}(\text{CH}_3)_2$  on 13.22 g. of Ga metal (99.99% pure) in the presence of a trace of mercuric chloride. Reaction took place in a sealed glass pressure tube at  $130^\circ$  over a period of 9 days.<sup>7</sup> The  $\text{Hg}(\text{CH}_3)_2$  was prepared by the method of Marvel and Gould.<sup>8</sup> Three fractionations of the crude  $\text{Ga}(\text{CH}_3)_3$  from a vacuum trap at room temperature, through a trap at  $0^\circ$ , and into a trap at  $-78^\circ$  gave a product with a vapor pressure of 65.4 mm. at  $0^\circ$ . This compares well with the literature values of 66.6,<sup>7</sup> 64.5,<sup>9</sup> and 65.9 mm.<sup>1</sup> The yield was nearly quanti-

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